How to use the CD-ROM

Arrangement of Data

Each semiconductor is labelled by a number "x.y", where x gives the number of the *substance group* ("Elements of the IVth group", "III-V compounds" etc.) and y the number of the *substance* ("C", "Si", "Ge" etc.) within the group.

For each substance the individual data are listed in six property groups:

- **Crystal structure** (lattice structure / space group / modifications / high temperature and high pressure phases).
- Electronic properties (band structure / energies at symmetry points of the band structure / energy gaps (direct energy gap, indirect energy gap) / exciton energies / intra conduction band energies / intra valence band energies / critical point energies / spin-orbit splitting energies / camel's back structure of the conduction band edge / structure of the top of the valence band / effective masses (electrons, holes) / g-factor of electrons / valence band parameters).
- Lattice properties (lattice parameters / linear thermal expansion coefficient / density / melting point / Debye temperature / heat capacity / phonon dispersion relations / phonon frequencies (wavenumbers) / sound velocities / second and third order elastic moduli / bulk modulus / Poisson ratio / internal strain parameter).
- Transport properties (electrical conductivity or resistivity (intrinsic conductivity) / (intrinsic) carrier concentration / carrier mobilities (electron mobility, hole mobility) / drift velocities and diffusion constants / thermal conductivity (resistivity) / Seebeck coefficient (thermoelectric power) / piezo- and elastoresistance coefficients).
- **Optical properties** (optical constants / absorption coefficient / reflectance / extinction coefficient / refractive index / dielectric function / dielectric constants / piezo- and elastooptic coefficients).
- **Impurities and defects** (binding energies of impurities / energy levels of impurities, defects and complexes or of deep centers).

If only few data are available some of the property groups are omitted or put together.

Location of substances

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To locate a substance you can open directly the bookmark "Data" and navigate via the bookmarks for the substance groups to the substances and from these to the property groups, references and figures.

Two bookmarks in the start program can help you to locate a substance you are interested in:

A. Navigation via substance groups

If you open the bookmark "Navigation via substance groups" a list of all semiconductors dealt with in this handbook is shown on the desktop ordered by 38 substance groups:

1 Elements of the IVth group and IV-IV compounds

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1.1 C (Diamond) 1.2 Si	1.3 Ge 1.4 Sn (grey Sn)	1.5 SiC 1.6 Si _x Ge _{1-x}	
2 III-V compounds			
2.1 BN	2.6 AlP	2.11 GaAs	2.16 InSb
2.2 BP	2.7 AlAs	2.12 GaSb	2.17 Al _{0.48} In _{0.52} As
2.3 BAs	2.8 AlSb	2.13 InN	2.17 Al _{0.49} In _{0.51} P
2.4 BSb	2.9 GaN	2.14 InP	2.17 AlAs _{0.96} P _{0.04}
2.5 AlN	2.10 GaP	2.15 InAs	
3 II-VI compounds			

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By clicking on the substance number the first data page of the respective substance is opened. In the bookmark list on the left side of the desktop the respective bookmark is marked, and you can easily open the sub-bookmarks for the property groups, references and figures.

B. Navigation via element systems

If you open the bookmark "Navigation via element systems" a list of all semiconductors dealt with in this handbook is shown on the desktop ordered alphabetically by the elements the substances are consisting of:

Ag–As–S	AgAsS ₂	28.1
	Ag ₃ AsS ₃	28.14
Ag–As–Se	AgAsSe ₂	28.2
Ag–As–Te	AgAsTe ₂	28.3

By clicking on the substance number the first data page of the respective substance is opened. In the bookmark list on the left side of the desktop the respective bookmark is marked, and you can easily open the sub-bookmarks for the property groups, references and figures.

1.3 Germanium, Ge

Crystal structure

The element germanium crystallizes under normal conditions in the diamond lattice (space group: cubic $O_h^7 - Fd3m$), see Fig. 1.0.2.

high pressure phases

Under pressure there is a transition from the diamond structure (Ge-I) to the tetragonal β -tin structure (Ge-II) around 10 GPa [63J, 78A, 84O, 86M]. Around 75 GPa there is a transformation from the β -Sn structure (Ge-II) to the simple hexagonal (sh) structure [86V], and around 100 GPa to an intermediate structure, and to the hexagonal close-packed structure [86V]. Recently it has been found that the β -tin transforms continuously to the hexagonal phase via an intermediate phase with Imma structure, as previously found in Si, predicted by [94L] and experimentally found by [96N]; this phase is not a mixture of the β -tin and hexagonal phases [96N]; a high-pressure double-hexagonal close-packed (dhcp) structure has been found above 100 GPa [86V].

Upon depressurising the Ge-II modification, Ge-III (ST12 structure, simple tetragonal with 12 nearest neighbours, space group P4₃2₁2, D₄⁸) is obtained at RT [63B, 64K, 65B, 83Q, 86M2], a mixture of Ge-III and Ge-IV at 170 ... 200 K, nearly pure Ge-IV (BC8 structure, body centred cubic with 8 nearest neighbours, space group Ia3, T_h^{7}) at 160 ... 170 K, a mixture of Ge-IV and a glassy phase at 130 ... 160 K, and mostly glassy Ge at 90 ... 130 K [92B].

Electronic properties

band structure : Fig. 1.0.10, Brillouin zone: Fig. 1.0.6

density of states in valence and conduction bands: Fig. 1.3.1.

The conduction band is characterized by eight equivalent minima at the end points L of the [111]-axes of the Brillouin zone (symmetry: L_6^-). The surfaces of constant energy are ellipsoids of revolution with their major axes along [111]. Higher energy minima of the conduction band are located at Γ and (above these) on the [100]-axes. The valence band is of the general structure shown in Fig. 1.0.7 with its maximum at Γ (symmetry: Γ_8^+), the light and heavy hole bands being degenerate at this point. Both bands are warped. The third spin-orbit split-off band has Γ_7^+ symmetry. In contrast to silicon the spin-orbit splitting energies are considerable. Thus, the symmetry notation of the double group of the diamond lattice is mostly used for Ge.

energies of symmetry points of the band structure (*E* relative to the top of the valence band)

$E(\Gamma_{6v})$	– 12.66 eV		theoretical data, Fig. 1.0.10a	76C
$E(\Gamma_{7v})$	– 0.29 eV		-	
$E(\Gamma_{8v})$	0.00 eV			
$E(\Gamma_{7c})$	0.90 eV			
$E(\Gamma_{6c})$	3.01 eV			
$E(\Gamma_{8c})$	3.22 eV			
$E(X_{5v})$	– 8.65 eV			
$E(X_{5v})$	– 3.29 eV		for experimental data from angular	
$E(X_{5c})$	1.16 eV		resolved photoemission, see Fig. 1.0.10	b
$E(L_{6v})$	– 10.39 eV		[85W] and [84H, 85N]	
$E(L_{6v})$	– 7.61 eV			
$E(L_{6v})$	– 1.63 eV			
$E(L_{4,5v})$	– 1.43 eV			
$E(L_{6c})$	0.76 eV			
indirect ene	ergy gap			
$E_{g,ind}(\Gamma_{8v}^{+}-$	L_{6c} 0.744(1) eV	T = 1.5 K	magnetotransmission	59Z

5,			•	
	0.664 eV	T = 291 K		
Eg,th	0.785 eV	T = 0 K (extrapol.)	temperature dependence of the	54M1
			intrinsic conductivity	

For $E_{g,ind}(T)$ below 425 K, see Fig. 1.3.2.

direct energy gap

Eg,dir	0.898(1) eV	T = 1.5 K
$(\Gamma_{8v}^{+} - \Gamma_{7c}^{-})$	0.805(1) eV	T = 293 K

temperature dependence of $E_{g,dir}$, see Fig. 1.3.2

exciton ground state

$E(1S_{3/2}^{3/2}(L_4^++L_5^+))$	+)) 740.46(3) meV	absorption (at 2.1 K) and luminescence	79M
$E(1S_{1/2}^{3/2}(L_6^+))$	741.58 (3) meV	(at 5.1 K), the energy of the assisting	
		LA(L)-phonon is subtracted, the exchange splitting has not been resolved so far	
Eb	4.18 meV	$1S_{3/2}^{3/2}(L_4^++L_5^+)$	76A
	3.17 meV	$1S_{1/2}^{3/2}(L_6^+)$	
$E(2S_{3/2}^{3/2})$	1.32 meV	theory including valence band	76A

 $E(2S_{1/2}^{3/2}) = 0.87 \text{ meV}$ $E(2P) = 2.35(5) \text{ meV} \quad T = 5 \text{ K}$ 3.13 (5) meV 3.35(5) meV 3.42(5) meV

critical point energies

E_1	2.111(3) eV	RT	from ellipsometric data by a two- dimen-	84V
			sional critical point analysis (2D CP)	
$E_1 + \Delta_1$	2.298(3) eV	RT	2D CP	
Δ_1	0.187(3) eV	RT	2D CP	
E'_0	3.123(19) eV	T = 100 K	3D CP	
$E'_0 + \Delta'_0$	3.309(19) eV	T = 100 K	3D CP	
$E'_{0,\mathrm{av}}$	3.110 eV	RT	2D CP; mean value of E'_0 and $E'_0 + \Delta'_0$	
E_2	4.368 (4) eV	RT	2D CP	
	4.346(3) eV	RT	1D CP	
spin-orbit spli	tting energies			
$\Delta_0(\Gamma_{7y}^+ - \Gamma_{8y}^+)$	0.297 eV	$T = 10 { m K}$	electroreflectance	75A

 $\begin{aligned} &\Delta_0(\Gamma_{7v}^{+} - \Gamma_{8v}^{+}) & 0.297 \text{ eV} \\ &\Delta'_0(\Gamma_{6c}^{-} - \Gamma_{8c}^{-}) & 0.200 \text{ eV} \\ &\Delta_1(\Lambda_{4,5v}^{-} - \Lambda_{6v}) & 0.184 \text{ eV} \\ &\Delta'_1(L_{4,5c}^{-} - L_{6c}^{-}) & 0.266 \text{ eV} \\ &\Delta''_1(L_{4,5v}^{-} - L_{6v}^{-}) & 0.228 \text{ eV} \end{aligned}$

conduction band, effective masses

$m_{n\perp}(L_6)$	$0.0823 m_0$	T = 120 K	magnetophonon resonance	82H
	$0.0807(8) m_0$	T = 30100 K	cyclotron resonance at 891 GHz	76F
$m_{n\parallel}(L_6)$	$1.59 m_0$	T = 120 K	magnetophonon resonance	82H
"	1.57 (3) <i>m</i> ₀	T = 30100 K	cyclotron resonance at 891 GHz	76F
$m_{\rm n}(\Gamma_7)$	$0.0380(5) m_0$	T = 30 K	piezomagnetoreflectance	70A

For the dependence of the transverse electron mass in the L_6 -minima on the energy above the bottom of the band, see Fig. 1.3.3a.

76B2,

76L

magnetoabsorption

degeneracy and conduction band

 $(2P_{3/2}^{5/2}, 2P_{3/2}^{3/2}), (2P_{-5/2}^{5/2}, 2P_{1/2}^{5/2})$

anisotropy

 $2P_{1/2}^{3/2}$

 $(2P_{3/2}^{3/2}, 2P_{3/2}^{5/2})$

 $(2P_{3/2}^{5/2}, 2P_{-5/2}^{5/2})$

far-infrared absorption

g-factor of electrons

<i>g</i> c	- 3.0(2)	T = 30 K		piezomagnetoabsorption	70A
valence band, ef	ffective masses				
m _{p,1}	$0.0438(3) m_0$ $0.0426(2) m_0$ $0.0430(3) m_0$ 0.2244(1)	T = 4 K, $T = 4 K,$ $T = 4 K,$	<i>B</i> [100 <i>B</i> [111 <i>B</i> [110	cyclotron resonance	56D
m _{p,h}	$\begin{array}{c} 0.284(1) \ m_0 \\ 0.376(1) \ m_0 \\ 0.352(4) \ m_0 \end{array}$	T = 4 K, T = 4 K, T = 4 K,	$B \parallel [100] B \parallel [111] B \parallel [110]$]]]	
m _{so}	$0.095(7) m_0$	T = 30 K		piezomagnetoabsorption	70A
For the depender	nce of the light hole m	hass on the e	energy bel	by the top of the band, see Fig. 1.3.3b.	
valence band pa	rameters				
A B C	- 13.3 - 8.57 12.78			extrapolated using a five level $k \cdot p$ scheme	75W2
Lattice propert	ties				
lattice paramete	er				
a	5.6579060 Å	T = 298.15	K	single crystal	75B
For the temperat	ure dependence of a in	n the range	20812°C	C, see Fig. 1.3.4.	
linear thermal e	expansion coefficient				
For temperature	dependence, see Fig.	1.3.5.			
α	0.00·10 ⁻⁶ K ⁻¹ 5.90·10 ⁻⁶ K ⁻¹ 7.20·10 ⁻⁶ K ⁻¹ 8.51·10 ⁻⁶ K ⁻¹	T = 0 K T = 300 K T = 600 K T = 1000 F	K	recommended values from different measurements	7582
density					
d	5.3234 g/cm ³	$T = 25^{\circ}\mathrm{C}$		hydrostatic weighing	52S
melting point					
$T_{\rm m}$ d $T_{\rm m}$ /d p	1210.4 K $- 2.04 \cdot 10^{-3}$ K atm ⁻¹				73H 67K
phonon dispersi	on curve, one phono	on density o	of states : 1	Fig. 1.3.6.	
phonon frequen	cies				
$v_{LTO}(\Gamma_{25'})$ $v_{TA}(X_3)$ $v_{LAO}(X_1)$ $v_{TO}(X_4)$ $v_{TA}(L_{3'})$ $v_{LA}(L_2)$ $v_{TO}(L_3)$	9.02(2) THz 2.38(2) THz 7.14(2) THz 8.17(3) THz 1.87(2) THz 6.63(4) THz 8.55(3) THz 7.27(2) THz	<i>T</i> = 300 K		coherent inelastic neutron scattering	72N

v_1	4.9138⋅10 ⁵ cm s ⁻¹	T = 298 K,	ultrasound measurement (20 MHz),	63M
v_2	$3.5424 \cdot 10^5 \text{ cm s}^{-1}$	$ ho = 45 \ \Omega \ \mathrm{cm}$	n-type samples	
v_3	5.3996·10 ⁵ cm s ⁻¹		designation: mode/direction of propaga-	
v_4	$3.5425 \cdot 10^5 \text{ cm s}^{-1}$		tion/direction of particle displacement	
υ ₅	2.7458·10 ⁵ cm s ⁻¹		v_1 : long/[001]/[001], v_2 : shear/[001]/[1 v_3 : long./[110]/[110], v_4 : shear/[110]/[v_4 : shear/[110]/[1	. 1 0] 001]
			05.shear/[110]/[110]	
elastic moduli				
<i>c</i> ₁₁	$12.40 \cdot 10^{11} \text{ dyn cm}^{-2}$	T = 298 K,	ultrasound measurement	71B2
<i>c</i> ₁₂	$4.13 \cdot 10^{11} \text{ dyn cm}^{-2}$	$ ho = 37 \ \Omega \ \mathrm{cm}$	temperature dependence: Figs. 1.3.7ac	
<i>c</i> ₄₄	$6.83 \cdot 10^{11} \text{ dyn cm}^{-2}$			
third order elas	stic moduli			
<i>c</i> ₁₁₁	$-7.10(6) \cdot 10^{12}$ dyn cr	m^{-2} $T = 298$ K,	from ultrasound measurements,	64M
<i>c</i> ₁₁₂	$-3.89(3) \cdot 10^{12}$ dyn cr	n^{-2} $\rho = 45 \Omega cm$	temperature dependence: Fig. 1.3.8	
<i>c</i> ₁₂₃	$-0.18(6) \cdot 10^{12} \text{ dyn cm}$	m ⁻²		
<i>c</i> ₁₄₄	$-0.23(16) \cdot 10^{12} \mathrm{dyn} \mathrm{c}$	cm ⁻²		
c ₁₆₆	$-2.92(8) \cdot 10^{12} \text{ dyn cr}$	m ⁻²		
c ₄₅₆	$-0.53(7) \cdot 10^{12} dyn cm$	m ⁻²		
Young's modul	us, torsion modulus :	see Fig. 1.3.9.		
bulk modulus				
Bs	$7.502 \cdot 10^{11} \mathrm{dyn} \mathrm{cm}^{-2}$	T = 298 K.	from ultrasound measurements	63M
5	5	$\rho = 45 \ \Omega \ \mathrm{cm}$		
internal-strain	parameter			
	0.76(4)		X-ray diffraction	64S
Debye tempera	ture			
$\Theta_{\rm D}(0)$	374 K		for $\Theta_{\rm D}(T)$, see Fig. 1.3.10	59F
heat capacity :	Fig. 1.3.11.			

Transport properties

sound velocities

Low field transport is maintained by electrons in the L_6^+ minima of the conduction band and holes near the point Γ_8^+ in the valence bands. At room temperature, the mobility of samples with impurity concentrations below 10^{15} cm⁻³ is limited essentially by lattice scattering; higher donor or acceptor concentrations result in an increasing influence of impurity scattering. At 77 K, even for doping concentrations below 10^{13} cm⁻³ the mobilities depend on the impurity concentration. Low temperatures and high concentrations lead to the replacement of free carrier conduction by impurity conduction.

intrinsic conductivity

σ_{i}	$2.1 \cdot 10^{-2} \Omega^{-1} \mathrm{cm}^{-1}$	T = 300 K	for temperature dependence, see Fig. 1.3.12	54M1
intrinsic carrier	concentration			

n _i	$2.33 \cdot 10^{13} \text{ cm}^{-3}$	T = 300 K	for temperature dependence,	54M1
			see Fig. 1.3.13	

The best fit to the curvein Fig. 1.3.13 is given by the empirical expression

 $n_i^2 = 3.10 \ 10^{32} \ T^3 \exp(-0.785/kT) \ cm^{-6} \ (kT \ in \ eV, T \ in \ K)$

$\mu_{ m n}$	3800 cm ² /Vs	<i>T</i> = 300 K	lattice mobility determined by conductivity measurements of high purity samples.	54M1
Temperature dep 1.3.14 and 1.3.1	pendence in the rang	ge 77300 K given by	$\mu_n = 4.90 \cdot 10^7 T^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^2/\text{Vs} (T \text{ in } \text{K}); \text{ see}^{-1.66} \text{ cm}^$	also Fig.
hole mobility				
$\mu_{ m p}$	1800 cm ² /Vs	T = 300 K	lattice mobility in high purity samples.	54M2
Temperature dep see also Fig. 1.3	pendence in the range.	ge 100300 K given b	by $\mu_{\rm p} = 1.05 \cdot 10^9 T^{-2.33} \text{ cm}^2/\text{Vs} (T \text{ in K});$	
piezoresistance	coefficients			
n-type material				
π_{11} π_{12} $(1/2)(\pi_{11}+\pi_{12}+$ $(1/2)(\pi_{11}+\pi_{12}-$	$-2.7 \cdot 10^{-12} \text{ cm}^2/\text{d}$ $-3.9 \cdot 10^{-12} \text{ cm}^2/\text{d}$ $\pi_{44}) -71.7 \cdot 10^{-12} \text{ c}$ $\pi_{44}) 62.0 \cdot 10^{-12} \text{ cm}$	yn $T = 300 \text{ K}$ yn $\rho = 5.7 \Omega \text{ cm}$ m ² /dyn ² /dyn	uniaxial tensile stress	54S
p-type material π_{11} π_{12} $(1/2)(\pi_{11}+\pi_{12}$	$\begin{array}{c} -3.7 \cdot 10^{-12} \text{ cm}^{2/3} \\ 3.2 \cdot 10^{-12} \text{ cm}^{2/3} \\ \pi_{44}) 48.1 \cdot 10^{-12} \text{ cm}^{-12} \end{array}$	dyn $T = 300$ K r dyn $\rho = 1.1 \Omega$ cm $r^{2/dyn}$	resistivity measurements	
elastoresistance	e coefficients			
n-type material				
m_{44} $(m_{11}-m_{12})/2$ $(m_{11}+2m_{12})/3$	- 93.0 0.4 - 6.6	T = 300 K, $\rho = 1.5 \Omega \text{ cm}$	calculated from measured piezoresistance coefficients with the aid of elastic constants tensor c_{ij}	54S
p-type material				
m_{44} $(m_{11}-m_{12})/2$ $(m_{11}+2m_{12})/3$	65.1 - 2.8 2.0	T = 300 K $\rho = 1.1 \Omega \text{ cm}$		
Seebeck coeffic	tient (thermoelectric	e power) (absolute valu	ues)	
S _n	$1.07 \cdot 10^{-3} \text{ V/K}$	T = 300 K, $\rho = 17.0 \Omega \text{ cm}$	temperature difference used: 1.54 K	65F
S _p	1.06·10 ⁻³ V/K	T = 280 K $\rho = 10.5 \Omega \text{ cm}$	temperature dependence: Fig. 1.3.16	

thermal conductivity

electron mobility

Temperature dependence: Fig. 1.3.17.

Optical properties

optical constants

real and imaginary parts of the dielectric constant measured by spectroscopical ellipsometry, *n*, *k*, *R*, *K* calculated from these data [83A2]. See also Fig. 1.3.18.

<i>hv</i> [eV]	ε_1	ϵ_2	n	k	R	$K [10^3 \mathrm{cm}^{-1}]$
1.5	21.560	2.772	4.653	0.298	0.419	45.30
2.0	30.361	10.427	5.588	0.933	0.495	189.12
2.5	13.153	20.695	4.340	2.384	0.492	604.15
3.0	12.065	17.514	4.082	2.145	0.463	652.25
3.5	9.052	21.442	4.020	2.677	0.502	946.01
4.0	4.123	26.056	3.90	3.336	0.556	1352.55
4.5	- 14.655	16.782	1.953	4.297	0.713	1960.14
5.0	-8.277	8.911	1.394	3.197	0.650	1620.15
5.5	- 6.176	7.842	1.380	2.842	0.598	1584.57
6.0	- 6.648	5.672	1.023	2.774	0.653	1686.84

refractive index , spectral dependence at RT

4.00541(11)	$\lambda = 8 \ \mu m$	mean values for ten samples from	82E1
4.00412(12)	9 µm	various suppliers measured at	
4.00319(11)	10 µm	20.0°C	
4.00248(10)	11 µm		
4.00194(11)	12 µm	temperature coefficient between	
4.00151(10)	13 µm	20°C and 25°C: 4.0·10 ⁻⁴ °C ⁻¹	

The spectral dependence of n, k, and K below 10 eV is shown in Fig. 1.3.19. For spectral dependence of n and k, see also Fig. 1.3.20. Temperature dependence: Fig. 1.3.21.

dielectric constant

п

$\mathcal{E}(\infty)$	16.00	T = 300) K	capacitance measurement	83S
d ln ε (∞)/dT	13.8(8) · 10 ^{−5} K ^{−1}	T = 77.	400 K	optical interference	59C
ε	16.5	T [K] 4.2	v [MHz] 750	capacitance bridge	66R
	16.0 16.2	4.2 300	9200 9200	microwave measurement	56A
	15.8	77	1	capacitance bridge	53D

For spectra of the real and imaginary parts of the dielectric constant, see Fig. 1.3.22. Temperature dependence of ε (∞): Fig. 1.3.23.

elastooptic constants

p_{11}	-0.154	T = 300 K,	interferometric technique	78F
p_{12}	-0.126	$\lambda = 10.6 \ \mu m$		
p_{44}	-0.073		polarimetric technique	

third order susceptibilities

The second order susceptibilities are zero as a result of inversion symmetry. The third order susceptibility has two independent components.

X 1111	$1.0(5) \cdot 10^{-10}$ esu	CO ₂ laser,	optical third order mixing	69W
X1122	$0.6(3) \cdot 10^{-10}$ esu	$\lambda = 10.6 \ \mu m$,		
χ_{1111}	$1 \cdot 10^{-10}$ esu	T = 300 K	calculated	70V

Impurities	and	defects
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shallow donors binding energies

binding energies of group V	substitutional	donors
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Impurity	$E_{\rm b} [{\rm meV}]$	<i>T</i> [K]	Experimental method, remarks	Ref.
Theory Sb Bi P As	9.81 10.29 12.81 12.89 14.17	10, 4 10 10 10	donor effective mass calculation optical absorption (photothermal ionization) optical absorption photoconductivity, absorption absorption	69F 75S1, 64R 64R 74S 64R
other shallow	w single dono	r binding ene	rgies	
Impurity	$E_{\rm b} [{\rm meV}]$	<i>T</i> [K]	Remarks	Ref.
Li O-Li O-H D $_1$ D $_2$ Ia Ib Ic	10.02 10.47 12.46 9.38 10.75 17.25 17.6 18.1	4 6 4	photoconductivity, absorption piezophotoconductivity photoconductivity photoconductivity, observed in oxygen doped crystals photoconductivity, observed in oxygen doped crystals photoconductivity, "thermal donor"-related photoconductivity, "thermal donor"-related photoconductivity, "thermal donor"-related	72S, 78H2 78H2 78H1, 80J 78C 78C 82C, 84C 82C, 84C 82C, 84C 82C, 84C
ground state	energies of g	roup VI subs	titutional donors	
Impurity	E [meV]	<i>T</i> [K]	Remarks	Ref.
Se ⁰	$E_{\rm c} - 268.2$	201	50 Hall effect, photoconductivity	59T, 85G 88G, 98O
Se ⁻	$E_{\rm c} - 512.4$	140	DLTS	
Te ⁰	$E_{\rm c} - 93$		Hall effect, photoconductivity,	57A, 59N, 59T, 88G
Te ⁻	$E_{\rm c} - 330$		DLTS	
S ⁰	$E_{\rm c} - 280$		Hall effect, DLTS	88G
S ⁻ O	$E_{\rm c} - 590$ $E_{\rm c} - 17$ $E_{\rm c} - 40$ $E_{\rm c} - 200$ $E_{\rm c} - 6080$		Hall effect, photoconductivity	61F, 62K, 63W, 82C 78E

$E_{\rm c} - 160180$			
$E_{\rm c} - 1617.3$	740	Absorption, D ⁰ state of oxygen-related	84C, 82C
		thermal double donor.	
$E_{\rm c} - 3437$	740	Singly ionized D ⁺ state of thermal donor	84C

shallow acceptors binding energies

binding energy of group III substitutional acceptors

Impurity	$E_{\rm b} [{\rm meV}]$	<i>T</i> [K[Remarks	Ref.
Theory B Al Ga	11.2 10.82 11.15 11.32	4, 8	acceptor EMT calculation photoconductivity, optical absorption full series of EMT like excited states	76B1 74H 65J
In Tl	11.99 13.45		optical absorption	65J

binding energy of other shallow singly charged acceptors

Impurity	$E_{\rm b} [{\rm meV}]$	Remarl	ζS	Ref.
H, Si(A ₂ /A ₁)	11.66, 10.59	hydrog 1.07 m strain d shift or	en-silicon associated quenched-in defect. eV ground state splitting due to lipole acting on central cell. + 21 μeV isotope a deuteration.	78H1, 85H1, 87K, 89D
H, $C(A_6)$	12.28, 11.30) 1.98 m	eV ground state tunnel splitting	78H1, 87K
CN	11.32	crystals	s grown in nitrogen atmosphere in	78H2, 85M, 85H1
substitutional	1	graphit	e coated crucible. 1.1 meV ground	
(A_3, A_5)		state sp	litting	
CN interstitial (A ₄)	10.77			
A ₇	11.01	no corr	elation with chemical impurity reported	78H1
SA ₁	8.69, 9.48	crystals	s quenched from 900°C. Detailed	81K, 81B, 82B
SA ₂	8.86, 9.59	identifi	cations uncertain. SA _{1,2} hole	
SA ₁ (Ni)	9.02, 9.21,	repulsi	ve cores since binding energy less	
	9.37, 9.62	than ef	fective mass value	
SA_1 (Cu)	9.73, 10.5			
SA ₂	13.89			
SA ₂	14.42			
SA ₃	17.89			
Be, H	11.29, 10.79	Be-dop 0.50 m	ed crystal grown in hydrogen atmosphere, eV ground state splitting	85M, 87K
Zn, H	12.53, 9.87	Zn-dop 2.66 m	eV ground state splitting	85M, 87K
A8	10.34			
A9	11.45 12.03	0.58 m	eV ground state splitting	82D
A10	11.43, 12.03	0.50 m	e v ground state spinning	02D
group II sub	stitutional d	ouble acce	ptors	
neutral doub	ble acceptor	binding en	ergies	
Impurity I	E _b [meV]	<i>T</i> [K]	Remarks	Ref.
Be^0 2	24.80	2,8	optical absorption and/or photoconductivity	67S, 71M, 74H, 83C
Zn ⁰ 3	32.98	2, 8	optical absorption and/or photoconductivity	60F, 71M
Mg^0 3	35.85		optical absorption and/or photoconductivity	85M
Cd^0 5	54.96		optical absorption and/or photoconductivity	71M
Hg ⁰ 9	91.65	620	optical absorption and/or photoconductivity	67C, 71M
Mn^0 5	55		optical absorption and/or photoconductivity	85M
singly ionize	d double acc	eptor bind	ling energies	
Impurity I	E _b [meV]	<i>T</i> [K]	Remarks	Ref.
Be ⁻	58.02		optical absorption	83C
Zn-	86.51	810	optical absorption, excited states observed	83C, 73B
Cd- 1	160	20	Hall effect, photoconductivity	59T
Hg ⁻ 2	230	10	optical absorption	67C
Mn ⁻ 1	100			85M

positively charged multiple acceptor binding energies

Impurity	$E_{\rm b} [{\rm meV}]$	<i>T</i> [K]	Remarks	Ref.
Be ⁺	5.1	1.2	photoconductivity	85M, 85L,
Zn^+	2.0	1.2	photoconductivity	83H
Mg^+	2.9	1.2	photoconductivity	
Mn ⁺	3.2	1.2	photoconductivity	
Hg^+	12.2	1.2	photoconductivity	
Cu^+	2.0	1.2	photoconductivity	
Ag^+	35	10	DLTS	89H

energy levels of group I substitutional triple acceptors and related complexes

energy levels of group I impurities

Impurity	$E [\mathrm{meV}]$	<i>T</i> [K]	Remarks	Ref.
Cu ⁰	$E_{\rm v} + 43.25$	20300	Hall effect, photo conductivity	54D, 57W, 85S1
Cu-	$E_{\rm v} + 330$	20300		57W, 84P1
Cu ²⁻	$E_{\rm c} - 260, E_{\rm v} + 410$	20300		57W, 84P1
Cu donor			Cu interstitial	64H
CuH ₂	$E_{\rm v}$ + 16.42, 16.81, 16.92,	10	photoconductivity (photothermal	77H1, 86K
	16.96, 17.03, 17.08,		ionization) produced by Cu diffusion	
	17.14, 17.21, 17.29,		in hydrogen atmosphere grown	
CuH	$E_{\rm w} + 175$		crystais	77H1
Cu. (D.H)	$E_{\rm v} + 18.10$			86K
Cu,(H,T)	$E_{\rm v} + 18.12$			86K
Cu, D ₂	$E_{\rm v} + 18.20$			86K
Cu, T ₂	$E_{\rm v} + 18.24$			86K
Cu, Li ₂	$E_{\rm v} + 20.41$		assignments of Cu, Li ₂	
			and Cu, (Li, H) may be reversed	85H2
Cu, (Li,H)	$E_{\rm v} + 25.30$			85H2
Cu, As	$E_v + 10.05/9.15$		split ground state, complex with	96S
			C_{3v} symmetry, found in Cu and	
			As doped, vacuum grown crystals	
Au donor	$E_{\rm V} + 44$		Hall effect, photoconductivity	55D, 57W,
Au ⁰	$E_{\rm v} + 135$		and DLTS, the DLTS data given	59N, 65O,
Au ⁻	$E_{\rm c} - 215$		here have not been corrrected for	87S
Au ²⁻	$E_{\rm c} - 56$		the capture cross section activation	
			energy. All levels are assigned to substitutional Au.	
Ag^0	$E_{\rm v} + 116$		Hall effect, DLTS	59N, 89H
Ag ⁻	$E_{\rm c} - 261$			
Ag ²⁻	$E_{\rm c} - 113$			

energy levels of transition metal impurities

Unless otherwise stated, transition metals give rise to acceptor levels. Ground state energy levels for the two ionization states (where relevant) are given.

Impurity	$E [\mathrm{meV}]$		T [K]		Remarks	Ref.
Cr	$E_{\rm v} + 70, E_{\rm v} +$	120			resistivity	59N
Mn	$E_{\rm v} + 160, E_{\rm c}$ -	- 370	304	00	resistivity, photoconductivity	55W, 59T
Fe	$E_{\rm v} + 350, E_{\rm c}$ -	- 270	304	00	resistivity, photoconductivity	54T, 59T
Co	$E_{\rm v} + 250, E_{\rm c}$ -	- 300	773	00	resistivity, Hall effect and	55T, 59T
					photoconductivity	
Co-donor	$E_{\rm v} + 83$				1 V	71B3
Ni ⁰	$E_{\rm v} + 207$				only level found with DLTS in	55T, 59T,
Ni - H	$E_{\rm v} + 159/E_{\rm v} +$	-299			hydrogen free, pure germanium the two levels correlate with the	92H, 92Z 92Z
					presence of both nickel and	
Pt	$E_{\rm v} + 40$				resistivity, three acceptor levels.	59N, 54D
Pd	$E_{\rm v} + 0.03, E_{\rm c}$	- 0.18			resistivity, Hall effect	80G
energy level	s of defect cente	rs			•	
Defect	E [meV]	Generatio	on	Rema	rks	Ref.
Di-vacancy				F + 1	100 120 160 meV levels are inter-	75M 77M
(V ₂)-donor			nretec	$E_V + 100, 120, 100 \text{ meV}$ levels are mer-		
(*2)-donor				and a	re double acceptors	
	$E_{\rm re} + 100$	a 1 MeV F			doned Ge	
	$E_{\rm V} + 100$ $E_{\rm v} + 120$	v irr at 280 K		Sh do	ned Ge	
	$E_{\rm V} + 120$ E_+160	γ-111. at 200 K		Bido	ped Ge	
Group V	$E_{\rm V} + 100$ E 200	1 MoV		n Ga	interpreted as structural modifi	64P 77M
interstitie1	$L_{\rm c} = 200$	1 IVIC V	NO V	n-Oc,	of group V interstitiel (D.)	04r, 7/101
mersunai		γ-111. at 2	.00 K	Cation	For group $\sqrt{100}$ mersular (D _i).	
				Annea bo int	als at 100 160°C. This level can also	
				irradi	ation	
Acceptors		v-irr at 4	2 К	(a) do	where acceptors appeal at 65 K	77M
labeled by		γ 111. at 4	2 К	interp	reted as vacancy-interstitial	//141
annealing				(Frenl	kel) pair "65 K" defects.	
8				(b) ac	ceptors, anneal at 160200 K, arise	
				from i	interstitial defects observed	
Di-vacancy	$E_{\rm v} + 80$			in dis	location free, high purity Ge	77H1, 77H2
(V ₂)–H	$E_{\rm v} + 200$			growr	n in H-atmosphere; also produced	82E2
				by γ-i	rr. of dislocated material.	
				Attrib	outed to di-vacancyhydrogen complexes	
				(c.f. d	li-vacancy-donor complexes above).	
				$E_{\rm v}$ + 7	71 meV obtained if degeneracy factor	
				of 4 a	ssumed.	
Di-vacancy	$E_{\rm v} + 100$	γ-irr. at 2	80 K	Li doj	ped, high purity Ge;	77H1
(V_2) –Li				resisti	ivity, Hall effect	
γ-irr. Ge	$E_{\rm v} + 270$	γ-irr. at 2	80 K	accep	tor, oxygen related defect	72C
γ-irr. Ge	$E_{\rm v} + 340$	γ-irr. at 2	80 K	accep	tor, silicon-self interstitial	77M
	F 00		00 37	intera	ction	0111
γ-irr. Ge	$E_{\rm c} - 80$	γ -irr. at 2	80 K	Hall e	ettect, Ge:Li	81V
γ-irr. Ge	$E_{\rm v}$ + 230, 380	γ-irr. at 2	80 K	DLTS	5	82P1
γ-irr. Ge	$E_{\rm c} - 420$					84P2

γ-irr. Ge	$E_{\rm c} - 200$	γ-irr. at 280 K	Sb doped Ge	81T
e ⁻ -irr. Ge	$E_{\rm c} - 200, 400$	e ⁻ -irr.	1.5 MeV irradiation, n-type Ge, DLTS studies	81F, 82F
e ⁻ -irr. Ge	$E_{\rm v}$ + 210, 240, 310	e ⁻ -irr.	10 MeV irradiation, p-type Ge, DLTS	83F1
e ⁻ -irr. Ge	$E_{\rm v}+240,290$	e ⁻ -irr.	n-type Ge (As doped)	
e [–] -irr. Ge	<i>E</i> _c – 260, 370, 380, 410	e ⁻ -irr.		82P2
		e ⁻ -irr.	3.5 MeV irradiation, DLTS studies	83A1
	<i>E</i> _c – 190, 260, 310, 430	e ⁻ -irr.	n-type Ge, 430meV trap dominant	83A1
	$E_{\rm v}$ + 160, 229, 350	e ⁻ -irr.	p-type Ge, 350 meV trap dominant	83A1
	2.4 μm infrared absorption band	35 MeV e ⁻ irr. 400 keV proton bombardment	vibrational levels of a complex defect. Attributed to di-vacancies in [75S1]. Anneals at 200 K. The electron irradiation gives rise to acceptor levels at E_v + 80 meV, E_v + 160 meV.	77G
			<110> axis from polarization experiments. Di-vacancy suggested.	82G
	2.72 um infra- red absorption band	fast irr.	IR absorption	75N
Oxygen	$E_{\rm c} - 0.13, 0.25,$	1.5 MeV	DLTS studies. Correlations with local	83F2
related defect	0.29	e ⁻ -irr.	mode spectra in [65W] for electron irradiated Ge:O suggested.	
	$E_{\rm v}$ + 0.27 eV	4 MeV	Resistivity	84L
	(donor) $E_{\rm v} + 0.14, 0.20$	e ⁻ -irr.		
	(acceptor) $E_{\rm c} - 200, 220, 270, 400$	thermal n	DLTS studies As, Sb doped Ge $E_{\rm c} - 220$ suggested to contain interstitial As	84F
Multi– vacancy complexes	$E_{v} + 16$ $E_{v} + 20$ $E_{v} + 14$ $E_{v} + 6080$ $E_{v} + 160180$	fast n irr.	acceptor levels obtained from Hall effect and resistivity measurements 17300 K. Observed in Sb-, As-, Ga- doped Ge. Attributed to multivacancy complexes – independent of impurity	77D, 55C, 70T
Edge dis-	$E_{\rm v} + 90$	plastically	Content Hall effect, resistivity	740, 75L,
location (?)	$E_{\rm v}$ + 75, 190, 270, 390	deformed Ge plastically deformed Ge	DLTS	79S 83B
Screw dislocation	$E_{\rm v} + 35,590$	twisted Ge	Hall effect, resistivity	75W1
Grownin dislocations			DLTS	85S2
	$E_{\rm v}+25,100$		p-type Ge	
	$E_{\rm c} - 90$		n-type Ge	
	$E_{\rm v} + 0.17, 0.36$	Laser irr. Ge	DLTS	83P1
"E _T I"	$E_{\rm c} - 40100$		DLTS on ultrapure Ge.	85H2

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Figures to 1.3

Fig. 1.0.2

The diamond lattice. The elementary cubes of the two face-centered cubic lattices are shown.



Fig. 1.0.6

Brillouin zone of the diamond lattice



Fig. 1.0.10

Ge. (a) Band structure obtained by a non-local pseudopotential calculation including spin-orbit interaction [76C], (b) comparison of the calculated valence bands along the Δ - and Λ -axes with angular resolved photoemission data [85W].



Ge. Density of states of valence and conduction bands (solid line: obtained by a non-local pseudopotential calculation [76C], dashed line: XPS spectra showing the valence band density of states [74L]). $E - E_v$ in eV.



Ge. Indirect and direct energy gaps vs. temperature [60M].



Ge. (a) Dependence of the transverse electron mass in the L_6 minima on energy above the bottom of the conduction band measured by magnetopiezotransmission [69A], (b) cyclotron mass of light holes vs. energy below the top of the valence band (B || [111], [110] and [100]) [80Z].



Ge. Lattice parameter vs. temperature [68S].



Ge. Linear thermal expansion coefficient vs. temperature. Experimental data from various authors and theoretical results (solid line) [85K].



Ge. Phonon dispersion curves (left panel) and phonon density of states (right panel) [91G]. Experimental data points [71N] and ab-initio calculations [91G].



Fig. 1.3.7a Ge. Elastic modulus c_{11} vs. temperature [53M].



Fig. 1.3.7b Ge. Elastic modulus c_{12} vs. temperature [53M].



Fig. 1.3.7c Ge. Elastic modulus c_{44} vs. temperature [53M].



Ge. Third-order elastic moduli vs. temperature. Solid lines: best fit to the data [83P2].



Ge. Young's moduli *E* and torsion modulus *G* vs. temperature according to [71B1] (*1*); [53F, 55F] (2); [53M, 59M] (3).



Ge. Debye temperature $\Theta_D(T)$ vs. temperature [59F].



Ge. Heat capacity at constant volume vs. temperature. Experimental points from [70T], solid line: theoretical [85K].



Ge. Conductivity vs. reciprocal temperature in the range of intrinsic conduction [54M1].



Ge. Intrinsic carrier concentration vs. reciprocal temperature [54M1].



Ge. (Ohmic) drift mobility of electrons vs. temperature obtained with a time-of-flight technique in hyperpure material (open circles); other symbols: data from five references; solid line: theory [81J].



Ge. Electron and hole mobilities vs. temperature for constant carrier concentration (high purity samples) [54M2].



Ge. Thermoelectric power times temperature vs. temperature. Full triangles: $n_a = 1.6 \cdot 10^{14} \text{ cm}^{-3}$; open circles: $n_a = 8.0 \cdot 10^{13} \text{ cm}^{-3}$; open triangles: $n_d = 8 \cdot 10^{12} \text{ cm}^{-3}$; full circles: $n_d = 8.9 \cdot 10^{13} \text{ cm}^{-3}$ [54G].



Ge. Thermal conductivity vs. temperature. (a) 3...400 K, (b) 400...1200 K. Solid curve in (a) and data in (b) from [64G], experimental data in (a) from [84B]. Dashed line in (b): extrapolated lattice component.



Ge. Real and imaginary parts of the dielectric constant vs. photon energy [83A2].



Ge. Spectral dependence of (a) the refractive index, (b) the extinction coefficient and (c) the absorption coefficient [59P].



Ge. Refractive index and extinction coefficient vs. photon energy in the region of the E_1 critical point at 120K [66P].



Ge. Temperature and wavelength dependence of the refractive index. Open circles Curve 1: $\lambda = 1.970 \mu m$; 2: $\lambda = 2.190 \mu m$; 3: $\lambda = 2.409 \mu m$; 4: $\lambda = 3.826 \mu m$; 5: $\lambda = 5.156 \mu m$ [60L].



Ge. Real and imaginary parts of the dielectric constant and energy loss function vs. photon energy [63P].



Ge. Temperature dependence of the high-frequency dielectric constant [96K]. Experimental data points [76I] and ab-initio calculations (full line); dotted line is the theoretical result without the effect of thermal expansion, and dashed line is the full line shifted to match the experimental data.

